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Influence of incommensurability on SDW and CDW amplitudes in underdoped cuprates

M.Eremin^a, I.Eremin,^a G.Seibold^b and S. Varlamov^b^aKazan State University,
420008 Kazan, Russian Federation^bCottbus Technical University,
03013 Cottbus, Germany

Self-consistent calculations of spin (charge) density wave order (SDW/CDW) parameters have been performed for bilayered cuprates on the basis of a singlet correlated band model. The smooth development of the pseudogap formation temperature is explained from underdoped to overdoped states and the Fourier amplitudes $\langle s_q \rangle$ (spin) and $\langle e_q \rangle$ (charge) modulations have been calculated. We have found a maximum of the incommensurability for doping $0.09 \div 0.11$ holes per copper site.

1. Basic relations

We start from the Hamiltonian:

$$H = \sum t_{ij} \Psi_i^{pd,\sigma} \Psi_j^{\sigma,dp} + \sum J_{ij} [2(s_i s_j) - \frac{n_i n_j}{2}] + \sum g_{ij} \delta_i \delta_j + H_{CDW} \quad (1)$$

where $\Psi_i^{pd,\sigma}$, $\Psi_j^{\sigma,dp}$ are quasiparticle Hubbard-like operators for the copper-oxygen singlet band, J_{ij} is the superexchange constant of the copper spin coupling and g_{ij} is a screened Coulomb repulsion of the doped holes, $1 + \delta_i = \sum_{\sigma} \Psi_i^{\sigma,\sigma} + 2\Psi_i^{pd,pd}$. The quasiparticle interaction H_{CDW} mediated by the phonon field leads to a CDW transition [1]. In addition to the usual mean field approach we have taken into account that anticommutators of Hubbard-like operators can be affected by the doping index level per one unit cell, spin magnetization and charge modulation (non-Fermi statistics effect (NFS)):

$$P_i^{\sigma} = \Psi_i^{pd,pd} + \Psi_i^{\sigma,\sigma} = \frac{1 + \delta_i}{2} + (-1)^{\frac{1}{2}-\sigma} s_i^z \quad (2)$$

The appearance of SDW and CDW we describe via the Fourier component:

$$\begin{aligned} s_{q_s}^z &= \frac{1}{N} \sum s_i^z \exp(iq_s R_i), \\ e_{q_e} &= \frac{1}{N} \sum \delta_i \exp(iq_e R_i) \end{aligned} \quad (3)$$

where q_s and q_e are the instability wave vectors with respect to SDW and CDW formation. Below we consider both as commensurate wave vectors $q_s = q_e = (\pi, \pi)$ and incommensurate wave vectors $q_s = q_e = (\pi \pm \varepsilon_x, \pi \pm \varepsilon_y)$. In the framework of a linear approximation we can write the equation for the order parameter:

$$\begin{aligned} \eta_{k+q}^{\sigma} &= [t_{k+q} - \frac{4}{(1 + \delta_0)^2} \langle s_i s_j \rangle t_k] \times \\ &< \frac{e_q}{2} - (-1)^{\frac{1}{2}-\sigma} s_q^z \rangle + G_{k+q}^{\sigma} + G_k^{ph} \end{aligned} \quad (4)$$

The CDW order parameters G_{k+q}^{σ} and G_k^{ph} are determined as in [2]. The thermodynamic values of the Fourier component $\langle e_q \rangle$ and $\langle s_q \rangle$ are calculated self-consistently:

$$\begin{aligned} e_q &= \frac{1}{2} \sum_k [\langle \Psi_{k+q}^{pd,\sigma} \Psi_k^{\sigma,pd} \rangle + \langle \Psi_{k+q}^{pd,\bar{\sigma}} \Psi_k^{\bar{\sigma},pd} \rangle] \\ s_q^z &= \frac{1}{2} \sum_k [\langle \Psi_{k+q}^{pd,\bar{\sigma}} \Psi_k^{\bar{\sigma},pd} \rangle - \langle \Psi_{k+q}^{pd,\sigma} \Psi_k^{\sigma,pd} \rangle] \end{aligned} \quad (5)$$